

Pairing Reentrance in warm rotating ^{104}Pd nucleus[†]

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The recent series of experiments conducted at the Bhabha Atomic Research Center (BARC) for the reaction $^{12}\text{C} + ^{93}\text{Nb} \rightarrow ^{105}\text{Ag}^* \rightarrow ^{104}\text{Pd}^* + p$ at the incident energy of 40 - 45 MeV has observed an anomalous enhancement of the nuclear level density (NLD) of ^{104}Pd nucleus at low excitation energy E^* and high angular momentum J^1 . This enhancement is similar to that previously predicted by the shell-model Monte Carlo (SMMC)² and FTBCS1 calculations³ for a warm rotating ^{72}Ge nucleus. Both the SMMC and FTBCS1 have pointed out that the local enhancement of NLD at low T and high J is associated with the pairing reentrance effect. The latter occurs when the angular momentum of the system is sufficiently high so that the pairing correlation, which is zero at low $T < T_1$, reappears at $T > T_1$. The goal of this work is to apply the FTBCS1 theory including finite angular momentum to study if the enhanced NLD observed in ^{104}Pd can be interpreted as the first evidence of pairing reentrance in a warm rotating finite nucleus.

The FTBCS1 theory at finite temperature and angular momentum is obtained based on the conventional finite-temperature Bardeen-Cooper-Schrieffer (FTBCS) theory that takes into account the effect of quasiparticle-number fluctuations (QNF) on the pairing field³. The numerical calculations are carried out for ^{104}Pd nucleus, whose single-particle spectra are taken from the axially deformed Woods-Saxon potential including the spin-orbit and Coulomb interactions. The quadrupole deformation parameter β_2 potential is adjusted so that the NLD obtained at different values of J fit best the experimental data, especially in the region where the enhancement of NLD is observed. The variation of β_2 with J is plotted in Fig. 1 (a). This figure clearly shows that ^{104}Pd nucleus undergoes a shape transition from the prolate shape ($\beta_2 > 0$) to the oblate one ($\beta_2 < 0$) at around $J = 20 \hbar$, which is reasonable in this mass region because of an alignment of protons in $g_{9/2}$ and neutrons in $h_{11/2}$ orbits. Figs. 1 (d) - (e) depict the NLD as a function of excitation energy E^* obtained within the FTBCS1 and the conventional FTBCS theories.

It is found that due to the QNF, the FTBCS1 gaps at different J values decrease monotonically with increasing E^* and do not collapse at the critical value $E^* = E_c^*$ as in the case of the FTBCS. As a result, the pairing reentrance takes place only in the pairing

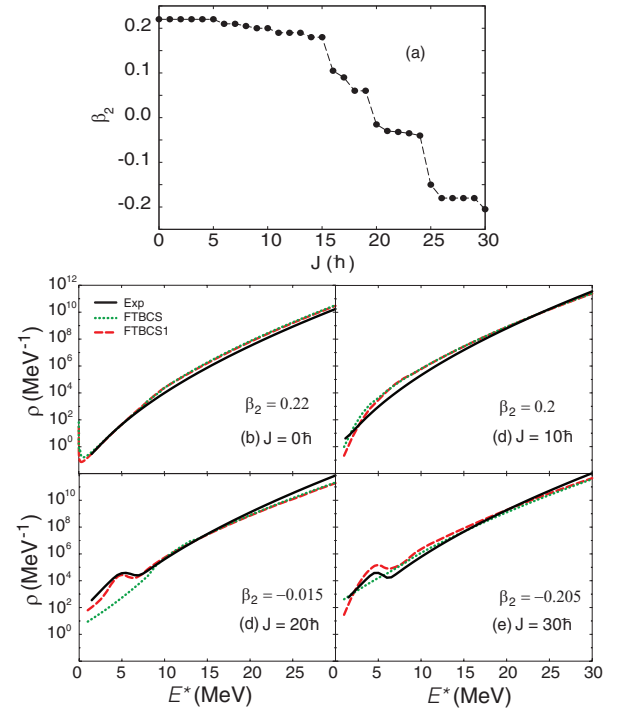


Fig. 1. (Color online) (a) - Quadrupole deformation parameter β_2 as functions of the total angular momentum J obtained within the FTBCS1 theory. [(b) - (e)] - Total NLD as function of excitation energy E^* obtained within the FTBCS (dotted lines) and FTBCS1 (dashed lines) at different values of J and β_2 . The solid lines are the experimental data.

gaps obtained within the FTBCS1 (*e.g.*, for protons at $J = 20 \hbar$ and neutrons and at $J = 30 \hbar$), whereas this effect does not appear in the FTBCS gaps. This leads to the local enhancements of the NLD obtained within the FTBCS1 at low E^* ($2 < E^* < 5$ MeV) and high J , in agreement with the experimental data. This agreement indicates that the observed enhancement of the NLD might be the first experimental detection of the pairing reentrance in a finite nucleus.

References

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