Long-range correlation of V⁰ particles in *p*-Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV with the ALICE detector

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Measuring correlations in particle production as a function of the azimuthal angular space and rapidity space is very useful for investigating particle production in high-energy nucleus-nucleus collisions. The long-range correlations in the rapidity space in nearside angular pairs of dihadron correlations were first observed in Au-Au collisions at $\sqrt{s_{NN}} = 200$ GeV at $\operatorname{RHIC}^{(1,2)}$. This long-range correlations are derived from the collective expansion of the initial geometry fluctuations. Unexpectedly, a similar structure has also been observed in high-multiplicity pp collisions at $\sqrt{s_{NN}} = 7$ TeV by the CMS experiment³). It is very interesting to study the correlation in p-Pb collisions because the initial gluon density and magnitude of the collective expansion are very different from those in other collision systems. The azimuthal anisotropy parameter v_2 of K, π , and p shows mass ordering at low transverse momentum $(p_{\rm T})$ and the trend is similar to Pb-Pb collisions⁴⁾. The mass ordering is a characteristic feature of collective expansion. This analysis aims to further explore the partonic collectivity by extracting v_2 of K_s^0 and Λ in p-Pb collisions at $\sqrt{s_{\rm NN}} = 5.02$ TeV with the ALICE detector. The correlations between the unidentified charged hadrons as trigger particle and K_s^0 and $\Lambda(\bar{\Lambda})$ as associated particles at $|\eta| < 0.8$ are measured as a function of the azimuthal angle difference $\Delta \phi$ and pseudo-rapidity difference $\Delta \eta$. K_s^0 and Λ decay into $\pi^+ + \pi^-$ and $p^+ + \pi^-$ with a characteristic decay pattern, called V⁰. Topological cuts are required to reduce the combinatorial background. The correlation function as a function of $\Delta \eta$ and $\Delta \phi$ between two charged particles is defined as $\frac{1}{N_{\text{trig}}} \frac{\mathrm{d}^2 N_{\text{asso}}}{\mathrm{d}\Delta \eta \mathrm{d}\Delta \phi} = \frac{S(\Delta \eta, \Delta \phi)}{B(\Delta \eta, \Delta \phi)}$, where N_{trig} is the total number of triggered particles in the event class and $p_{\rm T}$ interval, the signal distribution $S(\Delta\eta, \Delta\phi) = \frac{1}{N_{\rm trig}} \frac{{\rm d}^2 N_{\rm same}}{{\rm d}\Delta\eta {\rm d}\Delta\phi}$ is the associated yield per trigger particle from the same event, and the background distribution $B(\Delta\eta, \Delta\phi) = \alpha \frac{{\rm d}^2 N_{\rm mixed}}{{\rm d}\Delta\eta {\rm d}\Delta\phi}$ accounts for pair accounts and pair efficiency. for pair acceptance and pair efficiency. \vec{B} is constructed by taking the correlations between the trigger particles in one event and the associated particles in other events in the same event class. The α factor is chosen so that it is unity at $\Delta \eta \sim 0$ because the acceptance is flat along $\Delta \phi$. This correlation function is studied for different $p_{\rm T}$ intervals and different event classes. The correlation function in peripheral collisions is subtracted from that in central collisions to remove the auto-correlations from jets. Figure 1 shows the projection of the subtracted correlation functions onto $\Delta\phi$. To quantify azimuthal anisotropy (v_2) , the Fourier coefficients are extracted by fitting with the function $a_0 + 2a_1 \cos(\Delta\phi) + 2a_2 \cos(2\Delta\phi)$. The $v_{\rm n}$ coefficient can be obtained as $v_{\rm n}^{\rm K_s^0,\Lambda} = V_{\rm n}^{\rm K_s^0,\Lambda}/\sqrt{V_{\rm n}^{\rm h-h}}$, where $V_n^i = a_{\rm n}^i/(a_0^i + b^i)$, in which i is the index of h-h or h-V0 pairs (h denotes undefined hadrons) and b is the baseline determined by averaging over $1.2 < |\Delta\eta| < 1.6$ on the near side of the 60-100% event class. Figure 2 shows the extracted v_2 coefficient for K_s^0 and $\Lambda(\bar{\Lambda})$ compared to p and K as a function of $p_{\rm T}$. Mass ordering between the v_2 of K_s^0 and $\Lambda(\bar{\Lambda})$ as well as the kaon and proton is observed.

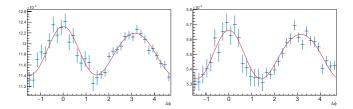


Fig. 1. Projection of the subtracted correlation functions of the associated K_s^0 (top) and Λ (bottom) yield per trigger particle with $1.5 < p_{T,trig}$, $p_{T,asso} < 2.5$ GeV.

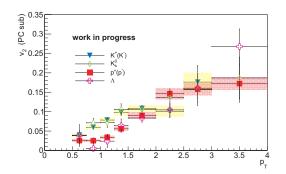


Fig. 2. v_2 of K_s^0 , $\Lambda(\bar{\Lambda})$ compared with one of kaon and proton. Error bars and shaded bands show statistical uncertainties and systematic, respectively.

References

- 1) STAR Collaboration, Phys. Rev. C80 064912 (2009)
- 2) PHOBOS Collaboration, Phys. Rev. Lett. 104 062301(2010)
- 3) CMS Collaboration, JHEP 09 091(2010)
- ALICE Collaboration, Phys. Lett. B, 726 164-177(2013)

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