# Long-range correlation of $\mathrm{V}^{0}$ particles in $p-\mathrm{Pb}$ collisions at $\sqrt{s_{N N}}=$ 5.02 TeV with the ALICE detector 

Y. Sekiguchi, ${ }^{* 1, * 2}$ H. Hamagaki,*2 and T. Gunji*2

Measuring correlations in particle production as a function of the azimuthal angular space and rapidity space is very useful for investigating particle production in high-energy nucleus-nucleus collisions. The long-range correlations in the rapidity space in nearside angular pairs of dihadron correlations were first observed in $\mathrm{Au}-\mathrm{Au}$ collisions at $\sqrt{s_{N N}}=200 \mathrm{GeV}$ at RHIC $^{1,2)}$. This long-range correlations are derived from the collective expansion of the initial geometry fluctuations. Unexpectedly, a similar structure has also been observed in high-multiplicity $p p$ collisions at $\sqrt{s_{N N}}=7 \mathrm{TeV}$ by the CMS experiment ${ }^{3)}$. It is very interesting to study the correlation in $p-\mathrm{Pb}$ collisions because the initial gluon density and magnitude of the collective expansion are very different from those in other collision systems. The azimuthal anisotropy parameter $v_{2}$ of $\mathrm{K}, \pi$, and p shows mass ordering at low transverse momentum ( $p_{\mathrm{T}}$ ) and the trend is similar to $\mathrm{Pb}-\mathrm{Pb}$ collisions ${ }^{4}$. The mass ordering is a characteristic feature of collective expansion. This analysis aims to further explore the partonic collectivity by extracting $v_{2}$ of $\mathrm{K}_{s}^{0}$ and $\Lambda$ in $p$ Pb collisions at $\sqrt{s_{\mathrm{NN}}}=5.02 \mathrm{TeV}$ with the ALICE detector. The correlations between the unidentified charged hadrons as trigger particle and $\mathrm{K}_{s}^{0}$ and $\Lambda(\bar{\Lambda})$ as associated particles at $|\eta|<0.8$ are measured as a function of the azimuthal angle difference $\Delta \phi$ and pseudo-rapidity difference $\Delta \eta . \mathrm{K}_{s}^{0}$ and $\Lambda$ decay into $\pi^{+}+\pi^{-}$and $p^{+}+\pi^{-}$with a characteristic decay pattern, called $\mathrm{V}^{0}$. Topological cuts are required to reduce the combinatorial background. The correlation function as a function of $\Delta \eta$ and $\Delta \phi$ between two charged particles is defined as $\frac{1}{N_{\text {trig }}} \frac{\mathrm{d}^{2} N_{\text {asso }}}{\mathrm{d} \Delta \eta \mathrm{d} \Delta \phi}=\frac{S(\Delta \eta, \Delta \phi)}{B(\Delta \eta, \Delta \phi)}$, where $N_{\text {trig }}$ is the total number of triggered particles in the event class and $p_{\mathrm{T}}$ interval, the signal distribution $S(\Delta \eta, \Delta \phi)=\frac{1}{N_{\text {trig }}} \frac{\mathrm{d}^{2} N_{\text {same }}}{\mathrm{d} \Delta \eta \mathrm{d} \Delta \phi}$ is the associated yield per trigger particle from the same event, and the background distribution $B(\Delta \eta, \Delta \phi)=\alpha \frac{\mathrm{d}^{2} \mathrm{~N}_{\text {mixed }}}{\mathrm{d} \Delta \eta \mathrm{d} \Delta \phi}$ accounts for pair acceptance and pair efficiency. B is constructed by taking the correlations between the trigger particles in one event and the associated particles in other events in the same event class. The $\alpha$ factor is chosen so that it is unity at $\Delta \eta \sim 0$ because the acceptance is flat along $\Delta \phi$. This correlation function is studied for different $p_{\mathrm{T}}$ intervals and different event classes. The correlation function in peripheral collisions is subtracted from that in central collisions to remove the auto-correlations from jets. Figure 1 shows the projec-

[^0]tion of the subtracted correlation functions onto $\Delta \phi$. To quantify azimuthal anisotropy $\left(v_{2}\right)$, the Fourier coefficients are extracted by fitting with the function $a_{0}+2 a_{1} \cos (\Delta \phi)+2 a_{2} \cos (2 \Delta \phi)$. The $v_{\mathrm{n}}$ coefficient can be obtained as $v_{\mathrm{n}}^{\mathrm{K}_{\mathrm{s}}^{0}, \Lambda}=V_{\mathrm{n}}^{\mathrm{K}_{\mathrm{s}}^{0}, \Lambda} / \sqrt{V_{\mathrm{n}}^{\mathrm{h}-\mathrm{h}}}$, where $V_{n}^{i}=a_{\mathrm{n}}^{\mathrm{i}} /\left(\mathrm{a}_{0}^{\mathrm{i}}+\mathrm{b}^{\mathrm{i}}\right)$, in which i is the index of $\mathrm{h}-\mathrm{h}$ or h -V0 pairs ( h denotes undefined hadrons) and b is the baseline determined by averaging over $1.2<|\Delta \eta|<1.6$ on the near side of the $60-100 \%$ event class. Figure 2 shows the extracted $v_{2}$ coefficient for $\mathrm{K}_{s}^{0}$ and $\Lambda(\bar{\Lambda})$ compared to p and K as a function of $p_{\mathrm{T}}$. Mass ordering between the $v_{2}$ of $\mathrm{K}_{s}^{0}$ and $\Lambda(\bar{\Lambda})$ as well as the kaon and proton is observed.


Fig. 1. Projection of the subtracted correlation functions of the associated $\mathrm{K}_{s}^{0}$ (top) and $\Lambda$ (bottom) yield per trigger particle with $1.5<p_{\mathrm{T}, \text { trig }}, p_{\mathrm{T}, \text { asso }}<2.5 \mathrm{GeV}$.


Fig. 2. $v_{2}$ of $\mathrm{K}_{s}^{0}, \Lambda(\bar{\Lambda})$ compared with one of kaon and proton. Error bars and shaded bands show statistical uncertainties and systematic, respectively.

## References

1) STAR Collaboration, Phys. Rev. C80 064912 (2009)
2) PHOBOS Collaboration, Phys. Rev. Lett. 104 062301(2010)
3) CMS Collaboration, JHEP 09 091(2010)
4) ALICE Collaboration, Phys. Lett. B, $\mathbf{7 2 6}$ 164177(2013)

[^0]:    *1 RIKEN Nishina Center
    *2 Center for Nuclear Study, the University of Tokyo

